

TWO-PHOTON IONIZATION OF THE HYDROGEN ATOM

L. P. RAPOPORT, B. A. ZON, and L. P. MANAKOV

Voronezh State University

Submitted August 5, 1968

Zh. Eksp. Teor. Fiz. 56, 400-401 (January, 1969)

An analytic expression in terms of the Appell function  $F_1$  is obtained for the two-photon ionization cross section of hydrogen atoms.

IN our previous work<sup>[1-3]</sup> analytic expressions were obtained for the amplitudes of two-photon processes in the hydrogen atom for transitions from one bound state to another. In the present note the proposed method is extended to the case of bound state-continuum transitions. As shown below, one can also obtain the amplitude for such transitions in analytic form.

For the sake of definiteness let us consider two-photon ionization of the ns-levels. The cross section of this process for the case of absorption of two photons of the same polarization and frequency  $\omega$  (Ry) has the following form in the nonrelativistic limit:

$$\frac{1}{I} \frac{d\sigma}{d\Omega} = \frac{4\pi r_0^2}{9a\omega^3 I_0} |M_0^{ns} + e^{i(hr-h_0)} M_2^{ns} (1 - 3 \cos^2 \theta)|^2. \quad (1)$$

Here  $\theta$  is the angle between the directions of polarization of the incident photons and the momentum  $\mathbf{p}$  of the outgoing electron,  $r_0$  is the classical radius of an electron,  $I$  is the intensity of the incident radiation in W/cm<sup>2</sup>,  $I_0 = 7.019 \times 10^{16}$  W/cm<sup>2</sup>. Let us consider then the case which is most important in practice, when  $\omega$  is less than the ionization potential of the atom. In this connection one can use the dipole approximation, and the reduced matrix elements of  $M_L$  have the form

$$M_L^{ns} = \int_0^\infty dr dr' (rr')^2 g_1(E_{ns} + \omega; r, r') \frac{dR_{ns}(r)}{dr} \times \left[ \frac{dR_L(r')}{dr'} + 3\delta_{L2} \frac{R_L(r')}{r'} \right], \quad (2)$$

where  $R_{ns}$  is the radial part of the wave function for the discrete spectrum,  $R_L$  and  $h_L$  are the radial part and phase of the L-th partial wave function for the continuous spectrum, having the asymptotic form of a plane wave plus a spherical incoming wave;<sup>[4]</sup>  $g_1$  is the radial part of the Coulomb Green's function for angular momentum  $l = 1$ ;  $\delta_{L2}$  is the Kronecker delta.

If the integral representation used in<sup>[1-3]</sup> is used for  $g_1$ , then all of the integrals appearing in Eq. (2) can be evaluated analytically. It is easy to see, by carrying out the integration in (2) first with respect to  $dr'$ , then with respect to  $dr$ , and finally with respect to the variable appearing in the integral representation, that one can represent  $M_L$  in the form

$$M_L^{ns} = N(\nu, k) P_{ns} \left( -\frac{\partial}{\partial x} \right) Q_L(x) |_{x=1/\nu}, \quad (3)$$

where

$$N(\nu, k) = \frac{32e^{-k\varphi}}{a(1 - e^{-2\pi k})^{1/2}(2 - \nu)\omega^3 \nu k^6}$$

$$Q_0 = -\frac{2}{(1 + x\nu)^4} \text{Im}[(2 - k^2 + 3ik)e^{i\varphi} F(0; 1; x)],$$

$$Q_2 = \frac{4k\nu[(1 + k^2)(4 + k^2)]^{1/2}}{(1 + x\nu)^4} \left[ F(0; 3; x) - \left( \frac{1 + x\nu}{1 - x\nu} \right)^2 \frac{2 - \nu}{4 - \nu} F(2; 3; x) \right],$$

$$F(s; t; x) \equiv F_1 \left( s + 2 - \nu; t - ik; 3 + ik; s + 3 - \nu; \frac{1 + x\nu}{1 - x\nu} e^{-i\varphi}, \frac{1 + x\nu}{1 - x\nu} e^{i\varphi} \right) \quad (4)$$

denotes the Appell function,<sup>[5]</sup>  $\varphi = 2 \tan^{-1}(\nu/k)$ ,  $\nu = n(1 - \omega n^2)^{-1/2}$ ,  $k = (ap)^{-1}$ ,  $a$  is the Bohr radius. The differential polynomial  $P_{ns}$  is obtained from

$$-a^{3/2} e^{r/an} \frac{d}{d(r/a)} R_{ns} \left( \frac{r}{a} \right)$$

by the substitution  $r/a \rightarrow -\partial/\partial x$ . In particular,

$$P_{1s} = 2, \quad P_{2s} = 2^{-1/2} \left( 1 + \frac{1}{4} \frac{\partial}{\partial x} \right).$$

The Appell functions entering into expressions (4) may be evaluated by means of expansion in rapidly converging series.<sup>[5]</sup> We have compared the cross section for ionization of the 2s-level with earlier calculations by Zernik.<sup>[6]</sup> Satisfactory agreement with<sup>[6]</sup> is obtained with the exception of regions of abrupt variation of the cross section ( $\omega = 0.137, 0.163, 0.188, 0.200$ ), where the results differ somewhat.

In conclusion we note that an examination of other types of bound state-continuum transitions may be carried out in similar fashion.

<sup>1</sup>B. A. Zon and L. P. Rapoport, ZhETF Pis. Red. 7, 70 (1968) [JETP Lett. 7, 52 (1968)].

<sup>2</sup>L. P. Rapoport and B. A. Zon, Phys. Lett. 26A, 564 (1968).

<sup>3</sup>B. A. Zon, N. L. Manakov, and L. P. Rapoport, Zh. Eksp. Teor. Fiz. 55, 924 (1968) [Sov. Phys.-JETP 28, 480 (1969)].

<sup>4</sup>H. A. Bethe and E. E. Salpeter, Quantum Mechanics of One- and Two-Electron Atoms, Academic Press, 1957 (Russ. transl., Fizmatgiz, 1960).

<sup>5</sup>A. Erdélyi, ed., Higher Transcendental Functions (Bateman Manuscript Project) (McGraw-Hill, 1953) Vol. 1 (Russ. transl., Izd. Nauka 1965).

<sup>6</sup>W. Zernik, Phys. Rev. 135, A51 (1964).

Translated by H. H. Nickle