ISOMER SHIFTS OF THE MÖSSBAUER Y LINE IN TRIVALENT ANTIMONY HALIDES

V. KOTHEKAR, B. Z. IOFA, S. I. SEMENOV, and V. S. SHPINEL'

Nuclear Physics Institute, Moscow State University

Submitted February 26, 1968

Zh. Eksp. Teor. Fiz. 55, 160-163 (July, 1968)

The isomer chemical shifts (ICS) of the Sb¹²¹ 37.2-keV Mössbauer γ line in trivalent antimony halides are measured. The value of $\Delta r/r$ (change of nuclear charge radius due to excitation) is obtained for Sb¹²¹ by comparing the ICS values for the tri- and pentavalent antimony compounds SbCl₃ and HSb HSbCl₆ · xH₂O and found to be (-9.5 ± 3) × 10⁻⁴. The ICS for SbHal₃ compounds increases linearly with increasing of ionicity of the bonds β . From a comparison of the observed slope of the straight-line plot of the ICS against β with the calculated value it follows that the change of the density $\psi_{\rm 5S}(O)^2$ due to screening by valence electrons in these compounds is insignificant.

In the investigation of the spectra of resonant absorption of 37.2-keV γ rays by Sb¹²¹ nuclei in various compounds of pentavalent antimony^[1], we obtained for the ratio $\rho = \Delta r/r$ of the changes in the radius of the distribution of the nuclear charge in Sb¹²¹ and Sn¹¹⁹ a value $\rho(\text{Sb}^{121})/\rho \text{Sn}^{119} = -5.5$, which agrees well with the results of Ruby et al.^[2]

We have continued the investigation of the isomer chemical shift (ICS) in halides of trivalent antimony $(SbCl_3, SbPr_3, SbI_3)$.

The spectra of resonant absorption of 37.2-keV γ rays by Sb¹²¹ nuclei were measured with an electrodynamic setup operating in constant velocity mode. The measurements were carried out at liquid-nitrogen temperature. The γ rays were recorded with a standard scintillation spectrometer with NaI (Tl) crystal (thickness 0.3 mm). The γ -radiation source was the isotope Sn¹²¹ in SnO₂.

In all the investigated compounds, the resonant spectrum was an unresolved asymmetrical line with noticeable broadening (see Fig. 1). Eight components of the quadrupole structure for the nuclear transition $7/2 \rightarrow 5/2$ could not be resolved, owing to the large width of the nuclear level ($\Gamma = 1.052$ mm/sec). The ICS obtained by us (δ) are listed in the table. The values of the ICS were obtained by adding eight peaks, the positions of which were determined by the equation

where

$$C(I, m) = [3m^2 - I(I+1)]/I(2I-1),$$

 $V_{m\bullet m} = \frac{c}{4E_{\chi}} e^2 q Q [RC(I^{\bullet}, m^{\bullet}) - C(I, m)] + \delta,$

R = Q*/q-ratio of the quadrupole moments of the excited and ground states, I*, I-spins and m*, m-their projections for the excited and ground states. The intensities of these peaks were determined by the squares of the Clebsch-Gordan coefficients^[3]. For an analysis of the experimental spectra, we used data on the nuclear quadrupole resonance (NQR)^[4,8] (see the table). According to the NQR data, the gradient of the electric field increases from iodine to chlorine, which is in good agreement with Γ_{exp} (experimental line width) for these compounds.

FIG. 1. Absorption spectrum of SBCl₃. The form of the spectrum corresponds to an unresolved hyperfine structure due to quadrupole interaction. The positions of the individual components of the hyperfine structure are shown by lines in the lower part of the figure. The velocity corresponding to the isomer shift is indicated by the arrow.



ANALYSIS OF THE EXPERIMENTAL DATA

The expression for the ICS, with allowance for the screening of the internal s-electrons and for the change of the density of the valence s-electrons on the Sb nucleus, owing to the change in the ionicity of the bond, can be written in the form

$$\delta = A \{ (1 - \beta_0) (N_s' - N_s) + 2b (N - N') N_s' \},$$
(2)

where

$$A = \frac{c}{E_{\gamma}} K(Z) \frac{\Delta r}{r} |\psi_{5s}(0)|_{\mathrm{Sb}^2},$$

 N'_{S} , N_{S} -number of valence s-electrons of the absorber and of the source, N', N-number of valence electrons, β_{0} -Crawford-Schawlow correction^[9], for which we assumed the value $\beta_{0} = 0.17$, the same as for Sn^{119 [10]}, and b-parameter that depends on the screening of the 5s electron by other electrons of the same shell.

X-ray diffraction studies of the compounds $\text{SbCl}_3^{[11]}$, SbBr₃^[12], and SbI₃^[13] show that the molecules of these compounds have a configuration of a trigonal pyramid with $\theta \approx 95.5^{\circ}$ (where θ -angle between the Hal-Sb-Hal

Compound	δ, mm/sec	n _{exp} mm/sec	$e^2 q Q_i \operatorname{Jn} \operatorname{Sb}^{(2)},$ MHz	e ² qQ on Hal., MHz	β
SbCl₃	-15.5 ± 0.2	4.3 ± 0.2	383,7 [³]	40,4 [⁵]	0,48
SbBr₃	-15.85 ± 0.2	3.35 ± 0.2	343,97 [⁴]	337 [⁵]	0,42
SbI₃	-16.5 ± 0.3	2.85 ± 0.2	169,74 [⁴]	1163 [⁶]	0,35

<u>Note</u>: β – ionicity of bonds calculated from the formula $\beta = 1 - Up - s^2$ (where s – hybridization of the bond ≈ 0.15 [⁷]).

(1)

Making up the hybrid wave functions for the antimony $atom^{[15]}$, we can obtain

$$N_s = 2 - 3\alpha - 3\alpha\beta, \quad N = 5 - 3\beta, \tag{3a}$$

where $\alpha = \cos \theta / (\cos \theta - 1) = 0.09$ gives the contribution of the s-electrons to the bonding orbital, and β is the ionicity of the bonds.

At the present time there is large scatter $(1.2-3.3)\times 10^{-4}$ in the values of $(\Delta r/r)\,Sn^{119}$ obtained by different authors $^{[10,16-19]}$. This uncertainty does not make it possible to obtain a reliable value of $\Delta r/r$ for Sn^{121} from the ratio $\rho(Sb^{121})/\rho(Sn^{119})$. Therefore, for an independent determination of $(\Delta r/r)Sb^{121}$, we compared the shifts for the compounds $SbCl_3$ and $HSbCl_6\cdot xH_2O$. For the second compound, on the basis of the sp^3d^2 hybrid functions, we can obtain

$$N_s = (1 - \beta), \quad N = 6(1 - \beta).$$
 (3b)

Using (2) and (3), we get for the ICS of these compounds

$$\delta_{(\text{SbCl}_3)} - \delta_{(\text{HSbCl}_4 \cdot \alpha \text{ H}_2 \text{ O})} = A \{ (1 - \beta_0) (1 - 3\alpha - 3\alpha\beta' + \beta) + 2b (1 - 6\beta + 3\beta') (2 - 3\alpha - 3\alpha\beta') \},$$
(4)

where β , β' -ionicities of the bonds in HSbCl₆ · xH₂O and SbCl₃. From the NQR data for SbCl₃ we obtain a value $\beta' = 0.48^{[19]}$. Starting from the electronegativities of the trivalent and pentavalent antimony (1.8 and 2.18)^[5], we can assume that β is 20% smaller. Substituting in (4) the measured values $\delta_{(SbCl_3)} = -15.5$ mm/sec and $\delta_{(HSbCl_6 \cdot xH_2O)} = -3$ mm/sec^[1] and neglecting the second term in the curly brackets (this results in an error < 2%) we obtain A = -15.1 mm/sec. Further, using K(Z) = -1.085 × 10⁻⁴ a_0^3 eV and the value $\psi_{5S}(O)^2$ = 18.28 a_0^{-3} (obtained in^[2] for the sp³ state of the antimony atom by the Hartree-Fock self-consistent field method), we obtain from Eq. (2)

$$\Delta r / r = (-9.5 \pm 3) \cdot 10^{-4}$$
.

This value of $\Delta r/r$ is 50% smaller than that obtained in^[1] under the assumption that the differences of the electronegativities $\Delta \chi = 2$ correspond to a 50% ionic bond, and for $\psi_{5S}(O)^2$ we used the nonrelativistic value 1.34×10^{26} cm⁻³. This value of $\Delta r/r$ is 10% larger than the results obtained by Ruby et al.^[2] This difference is apparently connected with the fact that in the present paper we used the experimental values of β , whereas in^[2] they assumed the atomic configurations s⁰p³ and sp³ for the compounds SbF₆⁻¹ and SbF₂, respectively.

We also obtained the dependence of δ on β for the compounds SbHal₃ (Fig. 2). From (2) and (3) we get for the ICS of compounds of the type SbHal₃

$$\delta = A \{ (1 - \beta_0) [2 - 3\alpha - 3\alpha\beta - N_s] + 2b[N - 5 + 3\beta] (2 - 3\alpha - 3\alpha\beta) \}.$$
(5)

Neglecting the term $3 \alpha \beta$, which is sufficiently small

FIG. 2. Dependence of the ICS for 37.2-keV gamma transitions in the Sb¹²¹ nucleus on the ionicity β of the bonds for the compounds SbCl₃, SbBr₃, and SbI₃.



compared with 2, we obtain for the slope of the dependence of the ICS on β the relation

$$\delta E / \delta \beta = A \{ (1 - \beta_0) (-3\alpha) + \delta b (2 - 3\alpha) \}.$$
(6)

Comparing (6) with the experimental slope ($\delta E / \delta \beta_{exp} = 8 \text{ mm/sec}$), we can estimate the constant b ~ -0.03.

It follows from this value of b that the change of the density at the nucleus due to the change in the number of the valence p-electrons in these compounds is insignificant, i.e., the change of the electron density at the nucleus due to the screening by the valence p-electrons is compensated by the change of the density due to the $p_{1/2}$ electrons.

¹V. A. Bryukhanov, B. Z. Iofa, V. Kothekar, S. I. Semenov, and V. S. Shpinel', Zh. Eksp. Teor. Fiz. 53, 1582 (1967) [Sov. Phys.-JETP 26, 912 (1968)].

²S. L. Ruby, G. M. Kalvius, G. B. Beard, and R. E. Snyder, Phys. Rev. **159**, 239 (1967).

³A. Simon, in: Deformatsiya atomnykh yader (Deformation of Atomic Nuclei) [Translation Collection], IIL, 1958, p. 365.

⁴R. G. Barne and P. J. Brays, J. Chem. Phys. 23, 1177 (1955).

⁵M. Haissinsky, J. phys. radium 7, 7 (1946).

⁶A. L. Schawlow, J. Chem. Phys. 22, 1211 (1954).

⁷T. P. Das and E. L. Hann, Nuclear Quadrupole

Resonance, Spectroscopy, Solid State Physics, 2d Supplement, 1958, p. 112.

⁸S. Ogawa, J. Phys. Soc. Japan 13, 618 (1958).

⁹M. F. Crawford and A. L. Schawlow, Phys. Rev. 76, 1310 (1949).

¹⁰T. P. Boquet, J. J. Chu, O. C. Kistner, M. L. Perlman, and G. T. Emery, Phys. Rev. Lett. **17**, 809 (1966).

¹¹J. Lindquist and A. Niggli, J. Inorganic and Nuclear Chem. 2, 345 (1956).

¹² J. Backer, Structural Reports 11, 272 (1947).

¹³J. Trotter and T. Zabel, Z. Krist. **128**, 67 (1966).

¹⁴ T. N. Polynova and M. A. Poraĭ-Koshits, Zh. strukt. khimii 7, 146 (1966).

¹⁵ W. Gordy, W. Smith and R. Trambarulo, Microwave Spectroscopy (Russ. Transl.), p. 280.

¹⁶ V. I. Gol'danskiĭ, G. M. Gorodinskiĭ, S. V. Koryagin et al., Dokl. Akad. Nauk SSSR 147, 127 (1962).

¹⁷A. J. F. Boyle, D. St. P. Bunbary, and C. Edwards, Proc. Phys. Soc. **79**, 416 (1962).

¹⁸T. Lee and P. A. Flinn, Phys. Lett. 19, 186 (1965). ¹⁹H. G. Robinson, H. G. Dehmelt and W. Gordy, J.

Chem. Phys. 22, 511 (1954).

Translated by J. G. Adashko

21