

POLARIZATION OF  $\Lambda$  HYPERONS PRODUCED IN  $\pi^-$ -C INTERACTIONS AT 7 GeV

V. A. BELYAKOV, WANG YUNG-CHANG, V. I. VEKSLER, N. M. VIRYASOV, TU YUAN-T'SAI, KIM HI IN, E. M. KLADNITSKAYA, A. A. KUZNETSOV, NGUYEN DIN TU, V. I. PENEV, and M. I. SOLOV'EV

Joint Institute for Nuclear Research

Submitted to JETP editor March 13, 1963

J. Exptl. Theoret. Phys. (U.S.S.R.) **45**, 90-92 (August, 1963)

A search was made for  $\Lambda$  hyperons produced in collisions between 7 GeV negative pions and carbon nuclei. The  $\Lambda$  hyperons are produced without longitudinal polarization.

THIS work was begun in 1960 to study the longitudinal polarization of  $\Lambda$  hyperons produced in complex nuclei. The longitudinal polarization is manifested in a forward-back asymmetry of the angular distribution of protons (or  $\pi^-$  mesons) from the  $\Lambda$  decay relative to the direction of the  $\Lambda$  hyperon momentum in the system in which the  $\Lambda$  hyperon is at rest. Up to the time our experiment was begun, only a few results had been published, and these were in disagreement with one another. Thus, in the study of the  $\Lambda$ -hyperon production by cosmic rays in complex nuclei<sup>[1]</sup> a very large forward-back asymmetry was observed, while no asymmetry was found in the production of  $\Lambda$  hyperons by negative pions at 1.12 and 1.23 GeV/c on hydrogen<sup>[2,3]</sup>. The problem was of definite interest, since the presence or absence of a longitudinal polarization is evidence of the violation or nonviolation of spatial parity conservation in strong interactions.<sup>[4,5]</sup>

It was very important to verify whether the asymmetry in the angular distribution observed in<sup>[1]</sup> depended on the energy of the incident particle or whether the interaction involved complex nuclei. Therefore, along with the study of the polarization of  $\Lambda$  hyperons produced in  $\pi^-$ -p interactions by 7-8 GeV  $\pi^-$  mesons,<sup>[6,7]</sup> we studied the polarization of  $\Lambda$  hyperons produced in  $\pi^-$ -C interactions by 7-GeV  $\pi^-$  mesons. For this purpose, we used the same pictures obtained with the 24-liter propane bubble chamber of the High-energy Laboratory, which was also used for the study of  $\Lambda$ -hyperon polarization in  $\pi^-$ -p interactions. All events with  $\Lambda$  hyperons and  $K^0$  mesons whose characteristics did not satisfy the selection criteria for  $\pi^-$ -p interactions<sup>[6,7]</sup> were considered to be  $\pi^-$ -C interactions.

It follows from<sup>[7]</sup> that of all the corrections the most important is the correction for the loss of the events in which the  $\pi^-$  meson from the  $\Lambda$  decay has a very short range. In our propane

chamber, such cases will be  $\Lambda$  decays with  $\pi^-$  mesons of range between 0 and 2 cm. Hence at the very outset of the scanning for  $V^0$  events, we paid special attention to such cases.

All  $V^0$  events whose kinematic characteristics corresponded to  $\Lambda$ -hyperon and  $K^0$ -meson decays were considered to be  $\Lambda$  hyperons, since it was shown in<sup>[7]</sup> that at our energies more than 90% of the  $V^0$  events whose decay kinematics are ambiguous are  $\Lambda$  hyperons.

It is known that  $\Lambda$  hyperons have a spin  $1/2$  and hence the angular distribution of the protons (and also the  $\pi^-$  mesons) from the  $\Lambda$  decay can be written in the form

$$f(\theta) = 1 + \alpha \bar{p} \cos \theta,$$

where  $\alpha$  is the asymmetry parameter,  $\bar{p}$  is the polarization component,  $\cos \theta$  characterizes the emission angle of the proton (or  $\pi^-$  meson) from the  $\Lambda$  decay in the  $\Lambda$  rest system (the direction of one of the coordinate axes coincides with the direction of the  $\Lambda$  momentum), where

$$\alpha \bar{p} = \frac{3}{N} \sum \cos \theta \pm \sqrt{\frac{3}{N} [1 - (\alpha \bar{p})^2]}$$

( $N$  is the total number of analyzed  $\Lambda$  hyperons).

For  $\alpha \bar{p}$  not equal to zero, the polarization will have a nonzero value, since  $\alpha \neq 0$  (thus,  $\alpha_\Lambda = 0.62 \pm 0.07$  according to Cronin<sup>[8]</sup> while  $\alpha_\Lambda = 0.68 \pm 0.07$  and  $\alpha_\Lambda = 0.66 \pm 0.20$  according to the data of Crawford and Leutner—Geneva Conference, 1962).

In our experiment, we found for 260  $\Lambda$  hyperons that

$$\alpha \bar{p}_1 = -0.01 \pm 0.11 \text{ for the forward back distribution,}$$

$$\alpha \bar{p}_2 = -0.06 \pm 0.11 \text{ for the right-left distribution,}$$

$$\alpha \bar{p}_3 = +0.04 \pm 0.11 \text{ for the up-down distribution.}$$

As before,<sup>[7]</sup> we paid special attention to the dependence of the quantity  $\alpha\bar{p}_1$ : 1) on the Λ-hyperon momentum  $p_\Lambda^*$  in the pion-nucleon c.m.s. and 2) on the multiplicity  $n_s$  of charged particles produced together with the Λ hyperons. We found

- 1)  $\alpha\bar{p}_1 = -0.02 \pm 0.14$  for  $p_\Lambda^* < 1000$  MeV,  
 $\alpha\bar{p}_1 = -0.02 \pm 0.16$  for  $p_\Lambda^* > 1000$  MeV,
- 2)  $\alpha\bar{p}_1 = -0.24 \pm 0.15$  for  $n_s \leq 3$ ,  
 $\alpha\bar{p}_1 = +0.24 \pm 0.15$  for  $n_s \geq 4$ .

We note that for Λ hyperons produced in π<sup>-</sup>-p interactions we found<sup>[7]</sup> the following values of  $\alpha\bar{p}_1$

- $$\alpha\bar{p}_1 = +0.15 \pm 0.10 \text{ for } n_s = 0 \text{ and } 2,$$
- $$\alpha\bar{p}_1 = -0.25 \pm 0.15 \text{ for } n_s = 4 \text{ and } 6.$$

The change in the sign of the quantity  $\alpha\bar{p}_1$  as a function of the charged-particle multiplicity in π<sup>-</sup>-p and π<sup>-</sup>-C interactions can be ascribed only to statistical fluctuations.

Thus the obtained results are in agreement with spatial parity conservation in strong interactions involving strange particles in interactions of 7-GeV negative pions with carbon nuclei.

The authors thank the large group of colleagues for discussions and valuable comments, the labora-

tory and technical staff for aid in the work, and the staff of the computing center for carrying out the calculations.

<sup>1</sup> R. A. Salmeron and A. Zichichi, *Nuovo cimento* **11**, 461 (1959).

<sup>2</sup> Crawford, Cresti, Good, Solmitz, and Stevenson, *Phys. Rev. Lett.* **1**, 209 (1958); *ibid.* **2**, 11 (1959).

<sup>3</sup> Lander, Powell, and White, *Phys. Rev. Lett.* **3**, 236 (1959).

<sup>4</sup> V. G. Solov'ev, *Nuclear Phys.* **6**, 618 (1958); *DAN SSSR* **129**, 68 (1959), *Soviet Phys. Doklady* **4**, 1255 (1959).

<sup>5</sup> A. Pais, *Phys. Rev. Lett.* **1**, 418 (1958).

<sup>6</sup> M. I. Solov'ev, *Proc. of the 1960 Ann. Intern. Conf. on High Energy Physics at Rochester, Univ. of Rochester, 1960*, p. 388.

<sup>7</sup> Veksler, Viryasov, Vrana, Kim, Kladnitskaya, Kuznetsov, Nguyen, Solov'ev, Hofmohl, and Ch'eng, *JETP* **44**, 84 (1963), *Soviet Phys. JETP* **17**, 58 (1963).

<sup>8</sup> J. W. Cronin, *Proc. of the 1962 Ann. Intern. Conf. on High Energy Physics at CERN, Geneva, 1962*, p. 453.

Translated by E. Marquit