

DECAY OF NUCLEI WITH LESS THAN 126 NEUTRONS

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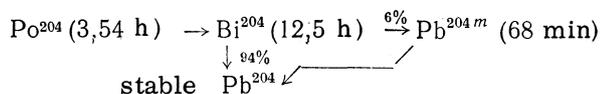
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The ratio of the  $\alpha$ -decay probability to the overall decay ( $\alpha$  decay plus E capture) probability  $\alpha/(\alpha + E)$  was measured for  $Po^{204}$ . The coefficients a and b in the linear dependence (1) for polonium isotopes with  $N < 126$  were determined. The values of  $\alpha/(\alpha + E)$  for  $Po^{196,198,200}$  were estimated with the aid of this linear relation. The values of the reduced derivative widths  $\delta_L^2$ , which characterize the probability of  $\alpha$ -particle formation were calculated for these isotopes from the value of  $T_{1/2}(\alpha)$ . The hindrance coefficients for a number of  $\alpha$  transitions were found for even-odd Po isotopes.

In previously published studies [1,2] of  $Po^{204}$  the ratio  $\alpha/(\alpha + E)$  of the probability of  $\alpha$  decay with an  $\alpha$ -particle energy  $E_\alpha = 5.37$  MeV to the overall decay probability ( $\alpha$  decay and E capture) was not measured directly. Karraker and Templeton [1] estimated this ratio to be  $\sim 1\%$  and Latimer et al. obtained the value  $(0.63 \pm 0.16)\%$  by interpolation.

We carried out an experiment in which the ratio of the probabilities  $\alpha/(\alpha + E)$  for  $Po^{204}$  was measured. To do this we compared the activity of  $Po^{204}$ , which undergoes  $\alpha$  decay and E capture, with the corresponding activity of  $Po^{206}$ , for which the  $\alpha$ -decay and E-capture rates are known [ $\alpha/(\alpha + E) = (5 \pm 1)\%$ ]. The amount of  $Po^{204}$  nuclei undergoing  $\alpha$  decay was measured relative to the  $\alpha$  decay of  $Po^{206}$  by means of an ionization  $\alpha$  spectrometer. The fraction of  $Po^{204}$  undergoing E capture was determined from the amount of Bi daughter nuclei separated chemically from Po. The Bi daughter-nucleus activity was measured with a  $4\pi$  scintillation counter under the assumption that the  $Po^{204}$  decay via E capture proceeds through the following radioactive chain [3]:



Here the  $Po^{204}$  yield in the Bi fraction was measured relative to the  $Po^{205}$  yield, while the  $Po^{205}$  yield was measured relative to the  $Po^{206}$  yield. It was then possible to determine the amount of nuclei undergoing E capture relative to the total amount of  $Po^{206}$ . A more detailed description of the experimental method can be found in [4].

On the basis of several runs, the value of the ratio  $\alpha/(\alpha + E)$  for  $Po^{204}$  was found to be

$(0.645 \pm 0.084)\%$ . The error represents the rms error for the series of measurements and does not include the error in the ratio of  $\alpha/(\alpha + E)$  for  $Po^{206}$ .

Using the experimental data listed in Table I for even-even isotopes of  $Po^{204}$  and  $Po^{200,202,206,208}$  (see [4-6]), we can plot the logarithm of the partial half-life for ground-state transitions as a function of  $(Q_{eff})^{-1/2}$ , where  $Q_{eff}$  is the total  $\alpha$ -decay energy with allowance for the recoil-nucleus energy and the influence of the electron shell of the atom (see [7], p.156). This dependence can be represented in the linear form (see figure)

$$\log T_{1/2}(\alpha) = a/\sqrt{Q_{eff}} + b. \tag{1}$$

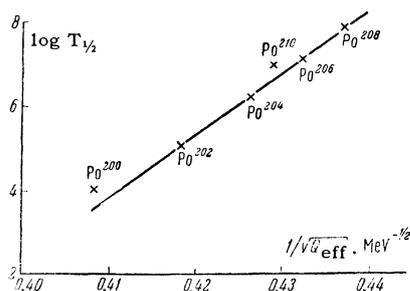
The position of the line was determined by the method of least squares from the data for  $Po^{202,204,206,208}$  for which the energy  $E_\alpha$  and the value of  $\alpha/(\alpha + E)$  have been more reliably determined.

If the partial half-life is measured in seconds and the total  $\alpha$ -decay energy is measured in MeV, then the coefficients for the series of even-even isotopes of Po ( $N < 126$ ) are:  $a = 150.36$  and  $b = -57.746$ . It has been shown earlier (see [7], p. 141) that, for even-even isotopes of Po with  $N > 126$ , these coefficients have the values  $a = 129.35$  and  $b = -49.9299$ . Hence the radioactive series of even-even isotopes of Po can be represented by linear dependences of type (1) with two sets of coefficients a and b. Here the isotopes with  $N < 126$  usually have an anomalously low probability for  $\alpha$  decay. The isotope  $Po^{210}$  does not lie on the line, owing to the effect of the shell with  $N = 126$  (see figure).

Using the dependence of  $\log T_{1/2}(\alpha)$  on Q, as determined by us we can calculate the values of  $T_{1/2}(\alpha)$  for  $Po^{196,198,200}$  and, from the experi-

Table I

Mass No., A	No. of neutrons	$E_\alpha$ , MeV	$T_{1/2}$	$\alpha/(\alpha+E)$ , %	$\delta_{L=0}^2$
210	126	5.299	138.4 d	110	0.027
208	124	5.108	2.93 y	~100	0.039
206	122	5.218	9.5 d	5	0.055
204	120	5.370	3.54 h	0.645	0.071
202	118	5.575	44.5 min	2	0.078
200	116	5.86	~10 min	~5( $16^{+4}_{-3}$ )	0.042 (0.13)
198	114	5.935	~6 min	( $23^{+7}_{-5}$ )	— (0.25)
196	112	6.14	~1.9 min	( $80 \pm 20$ )	— (0.19)



Dependence of logarithm of partial half-life  $T_{1/2}(\alpha)$  (in seconds) on total  $\alpha$ -particle energy for Po with  $N < 126$ .

mentally known overall half-lives  $T_{1/2}$ , we can then determine the ratios  $\alpha/(\alpha + E)$ . The results of the calculations are shown in Table I in parentheses. The discrepancy between the experimental and calculated values of  $\alpha/(\alpha + E)$  for  $\text{Po}^{200}$  can be ascribed to the experimental error which could occur in the case of an isotope with a short half-life and to the presence of  $\alpha$  particles from  $\text{Po}^{199}$  with almost the same energy ( $E_\alpha = 5.87$  MeV) [8].

Information on the partial half-life is of great interest from the viewpoint of  $\alpha$ -decay theory, in particular, in connection with the determination of the probability for  $\alpha$ -particle formation at the nuclear surface. The "reduced derivative width"  $\delta_{L=0}^2$ , i.e.,  $2\pi$  times the " $\alpha$  width" in the absence of the potential barrier [7] is frequently used as a characteristic of the  $\alpha$ -decay probability. The last column of Table I lists the values of  $\delta_{L=0}^2$  obtained with the aid of previously reported data [4,5,7] and the results of the present experiment. The values of  $\delta_{L=0}^2$  calculated with the aid of extrapolated values of  $T_{1/2}(\alpha)$  are shown in parentheses. For the calculation we used the rectangular nuclear potential model and assumed that the  $\alpha$ -particle angular momentum  $L$  was zero. The effective radius  $R$  was taken equal to  $9.3 \times 10^{-13}$  cm (see [7]). It is seen from the table that for even-even isotopes of Po with  $N < 126$  the value of  $\delta_{L=0}^2$  displays a tend-

ency to increase with the distance from the shell  $N = 126$ . Such a behavior has also been observed [7] for even-even isotopes of Po with  $N > 126$ . The value of  $\delta_{L=0}^2$  passes through a minimum in the region of the shell  $N = 126$ . Qualitatively, this behavior of  $\delta_{L=0}^2$  is in good agreement with the predictions based on the shell model [9].

For even-odd isotopes of Po with  $N < 126$ , we can calculate the forbiddenness coefficients  $F$  by means of the formula

$$\log F = \log T_{1/2}(\alpha) - (a/\sqrt{Q_{\text{eff}}} + b), \quad (2)$$

where the coefficients  $a$  and  $b$  are the same as those in the previously found logarithmic dependence. The calculated values of  $F$  are listed in Table II. The values of  $E_\alpha$ ,  $T_{1/2}$ , and  $\alpha/(\alpha + E)$  were taken from [4-6].

It is still difficult to draw general conclusions about the distribution of the hindrance coefficients, since experimental data on the level system for daughter nuclei and on the  $\alpha$ -particle angular momenta for  $\alpha$  transitions of Po are almost entirely lacking. It is seen from Table II that comparatively intense  $\alpha$  transitions for even-odd nuclei in the case of Po are, as a rule, weakly forbidden in comparison with  $\alpha$  transitions for even-even nuclei. For  $\text{Po}^{205}$  the anomalously small hindrance coefficient can be explained by the fact that it was actually calculated for a group of lines. This could

Table II

A	$E_\alpha$ , MeV	$T_{1/2}$	$\alpha/(\alpha+E)$ , %	F
209	4.877	100 y	~100	1.06
207	5.1	5.7 h	~0.01	2.04
205	5.2	1.5 h	0.074	0.32
203	5.48	47 min	~0.02	32
201	5.67	17.5 min	0.8	3.4
199	5.87	11 min	~7	2.8

also prove to be the case for other nuclei, since Forsling and Alväger established the presence of several groups of  $\alpha$  transitions in the case of some neutron-deficient Po isotopes. For  $\text{Po}^{203}$  the large hindrance coefficient can be related to the fact that this nucleus could have a neutron configuration with only one neutron in the unfilled  $f_{5/2}$  subshell. Then the  $\alpha$  particle is formed from nucleons of different subshells, and such an  $\alpha$  transition is, of course, forbidden. This suggestion is not in contradiction with the experimental value  $5/2$  for the spin of the  $\text{Po}^{203}$  nucleus<sup>[10]</sup>.

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<sup>3</sup>E. T. Hunter and J. M. Miller, Phys. Rev. **115**, 1053 (1959).

<sup>4</sup>Belyaev, Kalyamin, and Murin, Izv. An SSSR ser. fiz. **25**, 874 (1961), Columbia Tech. Transl. p. 886.

<sup>5</sup>Belyaev, Kalyamin, and Murin, Programma i tezisy dokladov XII Ezhegodnogo soveshchaniya po yadernoi spektroskopii (Program and Summaries of Reports of the 12th Annual Conf. on Nuclear Spectroscopy), Leningrad, 1962, AN SSSR, 1962, p. 65.

<sup>6</sup>Strominger, Hollander, and Seaborg, Revs. Modern Phys. **30**, 795 (1958).

<sup>7</sup>I. Perlman and J. Rasmussen, Alpha Radioactivity, Handb. d. Physik, Springer Verlag, Berlin, Gottingen, Heidelberg, 1957, Band 42, s. 109.

<sup>8</sup>W. Forsling and T. Alväger, Arkiv Fysik **19**, 353 (1961).

<sup>9</sup>H. J. Mang, Phys. Rev. **119**, 1069 (1960).

<sup>10</sup>S. Axensten and C. M. Olsmats, Arkiv Fysik **19**, 461 (1961).

Translated by E. Marquit