

$U^{238}$ , and  $Cf^{252}$ ). Secondly, the gamma quanta energies depends little on the excitation energy of the compound nucleus prior to fission.

The authors express their gratitude to Yu. I. Belyanin for insuring operation of the accelerated tube in the performance of this experiment.

<sup>1</sup>R. B. Leachman, Phys. Rev. **101**, 1005 (1956). Artem'ev, Protopopov, and Shiryaev, Труды РИАН (Trans. Radium Inst. Acad. Sci.) **9**, in press.

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<sup>3</sup>A. N. Protopopov and B. M. Shiryaev, J. Exptl. Theoret. Phys. (U.S.S.R.) **34**, 331 (1958), Soviet Phys. JETP **7**, 231 (1958).

Translated by J. G. Adashko  
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### ZEMPLEN'S THEOREM IN RELATIVISTIC HYDRODYNAMICS

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Submitted to JETP editor December 12, 1958

J. Exptl. Theoret. Phys. (U.S.S.R.) **36**, 956 (March, 1959)

KHALATNIKOV<sup>1</sup> has shown that for a relativistic shock wave of low intensity the theorem of Zemplen and the conditions of mechanical stability,  $v_1 > c_1$ ,  $v_2 < c_2$ , are applicable provided only that the following inequality holds;

$$\left(\frac{\partial^2(w/n)}{\partial p^2}\right)_s > 0 \quad (1)$$

(where  $w$  is the heat function per particle,  $s$  the entropy per particle,  $n$  the density of particles measured in the rest system of the particles, and  $p$  the pressure.)

These results are also applicable for relativistic shock waves of any intensity. The proof can be done in a similar way to Landau and Lifshitz, (reference 2, paragraph 84,) for the case when the shock adiabate lies in the plane ( $p, w/n$ ). In this case, formula (84,6) will correspond to

$$w_2 T_2 ds_2 = \frac{1}{2} (w_1/n_1 - w_2/n_2)^2 d(j^2),$$

and the expression

$$1 - \frac{v_2^2}{c_2^2} = (V_1 - V_2) \left[ 1 - \frac{j^2 (V_1 - V_2)}{2T_2} \left( \frac{\partial V_2}{\partial s_2} \right)_{p_2} \right] \frac{d(j^2)}{dp_2}$$

is replaced by

$$1 - \frac{u_2^2}{a_2^2} = \left( \frac{w_1}{n_1} - \frac{w_2}{n_2} \right) \left[ 1 - \frac{j^2 (w_1/n_1 - w_2/n_2)}{2w_2 T_2} \left( \frac{\partial (w_2/n_2)}{\partial s_2} \right)_{p_2} \right] \frac{d(j^2)}{dp_2},$$

$$j = nu, \quad u = v/\sqrt{1-v^2}, \quad a = c/\sqrt{1-c^2},$$

(where  $c$  is the velocity of sound, and the velocity of light is taken as unity.) It follows from this that the quantity  $n/w$ , as well as the pressure and the density, are increased on the shock wave.

The inequality (1), for the nonrelativistic case, reduces to the well known conditions,  $(\partial^2(1/n)\partial p^2)_s > 0$ . For a relativistic ideal gas we have

$$\left(\frac{\partial^2(w/n)}{\partial p^2}\right)_s = \frac{2(2-\gamma)}{\gamma(\gamma-1)} \frac{1}{pn^2}.$$

The last expression is always positive, since the quantity  $\gamma$  is within the interval<sup>3</sup>  $1 < \gamma \leq 5/3$ .

It should be noted that for an ultra-relativistic ideal gas,<sup>2</sup>  $\gamma = 4/3$ .

The author wishes to thank A. I. Akhiezer and G. Ya. Lyubarski for valuable discussions.

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<sup>2</sup>L. D. Landau and E. M. Lifshitz, Механика сплошных сред (Mechanics of Continuous Media), GITTL, M. 1954.

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Translated by S. Kotz  
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### ON ELECTROMAGNETIC SHOCK WAVES IN FERRITES

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Submitted to JETP editor December 18, 1958

J. Exptl. Theoret. Phys. (U.S.S.R.) **36**, 957-958 (March, 1959)

WE investigate the propagation of a uniform plane electromagnetic wave in a medium with non-linear dependence of the induction  $\mathbf{B}$  on the magnetic field  $\mathbf{H}$ .\* We assume to begin with that the